

1 **Title: Observational signatures of transuranic fission fragments in stars**

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36 **Abstract:**

37 The heaviest elements observed in the universe are produced by the rapid neutron-capture
38 process (r-process). Its termination among transuranic nuclei is poorly understood because these
39 elements are inaccessible to experiments, forcing nuclear models to extrapolate from limited
40 constraints. We show that the elements Ru, Rh, Pd, and Ag (atomic numbers $44 \leq Z \leq 47$, mass
41 numbers $99 \leq A \leq 110$) exhibit a correlation with abundances of heavier elements ($63 \leq Z \leq 78$,
42 $A > 150$) that is not shared by their immediate neighbors ($34 \leq Z \leq 42$ and $48 \leq Z \leq 62$) for stars that are
43 enhanced in r-process elements. Coproduction via fission fragments of transuranic nuclei
44 provides the most compelling explanation for this behavior. We conclude that this signature
45 provides the first evidence that neutron-rich fissioning nuclei with mass numbers > 260 are
46 produced in r-process events, such as neutron-star mergers.

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49 **One-Sentence Summary:**

50 A meta-analysis of stellar chemical abundances uncovers trends indicative of fission of the
51 heaviest elements.

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54 Understanding the origin of the elements is one of the major challenges of modern astrophysics.
55 Stars and stellar remnants synthesize the heaviest elements listed on the periodic table through
56 the rapid neutron-capture process, or r-process. The 2017 merger of a pair of neutron stars
57 detected in gravitational waves and electromagnetic radiation (1) established these types of
58 events as sites where the r-process occurs. The composition of material ejected constrains the
59 conditions present during the merger (2) and properties of the progenitor neutron stars, including
60 their equation of state (3). Freshly produced lanthanide elements ($57 \leq Z \leq 71$) (4, 5) and Sr ($Z =$
61 38) (6) were likely present in the ejecta of that 2017 event, but otherwise its detailed chemical
62 composition could not be inferred from either the light curve or spectra of its afterglow.

63
64 In contrast, the detailed compositions of a small number of ancient stars in the Milky Way can be
65 derived from hundreds of absorption features of more than 40 r-process elements detectable in
66 their spectra. The abundance patterns of heavy r-process elements found in these stars are
67 similar, which is often referred to as the "universality" of the r-process. The chemical inventory
68 of each star links directly to the ejecta of an individual r-process event, such as a neutron star
69 merger or a rare type of supernova that occurred in the early Universe and enriched the gas from
70 which these stars formed (7).

71
72 We assemble a sample of 42 stars in the Milky Way from data presented in 35 studies in the
73 literature (8). We select stars whose heavy elements were formed via the r-process, without
74 contamination from other processes, such as the slow neutron-capture process (s-process). We
75 follow common practice in adopting the [Fe/H] ratio (defined as $\log_{10}(N_{\text{Fe}}) - \log_{10}(N_{\text{Fe}})_{\odot}$, where
76 N is the number density and the \odot subscript indicates the Solar value) as a measure of the overall
77 metal enrichment of a star, and the [Eu/Fe] ratio (defined as $\log_{10}(N_{\text{Eu}}/N_{\text{Fe}}) - \log_{10}(N_{\text{Eu}}/N_{\text{Fe}})_{\odot}$) as

78 a measure of the enhancement of r-process elements relative to the overall metal enrichment. Our
79 sample spans $-3.57 \leq [\text{Fe}/\text{H}] \leq -0.99$ and $-0.52 \leq [\text{Eu}/\text{Fe}] \leq +1.69$ (Tab. S1–S4). Fig. 1
80 illustrates the heavy-element abundance patterns in these stars (see also Fig. S1 and S2). Stars
81 with higher $[\text{Eu}/\text{Fe}]$ ratios, which are marked with larger symbols, generally exhibit abundances
82 of some elements (including Ru, Rh, Pd, Ag, Gd, Tb, Dy, and Yb) that are slightly enhanced
83 relative to stars with lower $[\text{Eu}/\text{Fe}]$ ratios. This excess is not expected in the paradigm of r-
84 process universality, and we now aim to explain it.

85

86 We empirically calculate a baseline abundance pattern (Tab. S5 and S6) using the 13 stars (30%
87 of the sample) with the lowest levels of r-process enhancement, $[\text{Eu}/\text{Fe}] \leq +0.3$. Fig. 2 shows the
88 abundance excess, which is the difference between the individual chemical abundance
89 measurements in each star and the baseline abundance for that element. Three sets of elements
90 behave similarly in our sample: Se, Sr, Y, Zr, Nb, and Mo ($34 \leq Z \leq 42$); Cd, Sn, and Te ($48 \leq Z$
91 ≤ 52); and Ba, La, Ce, Pr, Nd, and Sm ($56 \leq Z \leq 62$). They are grouped together in panels A, C,
92 and D, respectively. Ratios among these elements exhibit no significant correlation with $[\text{Eu}/\text{Fe}]$,
93 and therefore no excess (8).

94

95 Two sets of elements exhibit significant positive correlations with $[\text{Eu}/\text{Fe}]$ (8): Ru, Rh, Pd, and
96 Ag ($44 \leq Z \leq 47$); and Gd, Tb, Dy, Ho, Er, Tm, Yb, Hf, Os, and Pt ($64 \leq Z \leq 78$; plus Eu, $Z =$
97 63). They are grouped together in panels B and E, respectively, of Fig. 2. These correlations
98 signify a newly identified extension of r-process universality that links many of the heavier r-
99 process elements with a few lighter ones. Previous work has shown that the range of elements
100 that participate in universality is bounded by two important deviations (supplementary text). One
101 is characterized by large (> 1.5 dex) differences in the overall amounts of lighter r-process

102 elements (atomic numbers $34 \leq Z < 56$) relative to the heavier ($Z \geq 56$) ones. The other is
103 characterized by small (< 0.7 dex) variations in the abundances of the actinide elements Th ($Z =$
104 90) and U ($Z = 92$) relative to other heavy r-process elements. Our findings also represent
105 deviations from standard two-component models of r-process nucleosynthesis (9, 10), where one
106 component (a weak, or limited, r-process) is mainly responsible for the lighter r-process elements
107 and another (a main r-process) is mainly responsible for the heavier r-process elements.

108

109 We propose that these element groups ($44 \leq Z \leq 47$ and $63 \leq Z \leq 78$) are coproduced as
110 transuranic ($Z > 92$) fission fragments in the r-process, establishing a natural connection between
111 them. Models predict that transuranic elements are produced in r-process events when the ejecta
112 contain very neutron-rich material (11–13). In these cases, the synthesis of heavy elements
113 terminates when transuranic nuclei undergo fission and repopulate the r-process chain at lower
114 masses (14). Fission deposition effectively "washes away" variations in the initial conditions
115 (15), leading to robust relative abundances among the fission products, such as Ag and Eu. (16)
116 recognized this behavior as a potential signature of fission, through a comparison of theoretical
117 calculations with Pd, Ag, and Eu abundances in a limited sample of stars. This behavior can be
118 seen in Fig. 3, which uses neutron star merger dynamical ejecta (8) to demonstrate the impact of
119 fissioning on the $\log \epsilon(\text{Ag}/\text{Eu})$ ratios. Models that do not include a fission component (panel A)
120 are sensitive to differences in initial conditions, such as progenitor mass, and they predict
121 varying $\log \epsilon(\text{Ag}/\text{Eu})$ ratios in the ejecta. Additionally, neutron star merger dynamical ejecta is
122 accompanied by ejecta from the post-merger accretion disk, which likely contains few or no
123 fissioning nuclei, so additional ejecta from the disk could lead to even larger variations from one
124 event to another. Models that include a fission component (panel B) predict that the mass-
125 weighted $\log \epsilon(\text{Ag}/\text{Eu})$ ratios are approximately constant, and they match the observed stellar

126 ratios. Similar behavior is expected for other candidate sites. For example, the neutron richness
127 predicted for the ejecta of magneto-rotationally driven supernovae is known to vary based on the
128 magnetic field of the progenitor (17). In any site capable of hosting fission, fission deposition
129 acts to stabilize the abundance ratios of coproduced elements against such astrophysical
130 variations.

131
132 Figure 1 illustrates the feasibility of this proposed scenario. The green bands in Fig. 1 mark the
133 baseline abundance pattern, which we calculate from the stars that we assume contain minimal
134 contributions from fission deposition. We add contributions from transuranic fission fragments,
135 enhanced by factors of 1, 2, and 4 relative to the baseline pattern (8). These enhancement factors
136 span the range of abundance excesses found for elements with $44 \leq Z \leq 47$ and $63 \leq Z \leq 78$ in the
137 stars with the highest levels of [Eu/Fe], which are marked by the largest symbols in Fig. 1.
138 Fission fragment deposition may contribute up to $\approx 75\%$ of the production of the elements with
139 the largest observed excesses, including Pd, Ag, Gd, and Yb (8).

140
141 Substantial contributions from fission fragments of transuranic nuclei provide the most
142 compelling explanation for the correlations between elements with $44 \leq Z \leq 47$ and $63 \leq Z \leq$
143 78. No other known nucleosynthesis process, including the weak or main s-process, intermediate
144 neutron-capture process, or weak r-process, can reproduce the observed behavior (supplementary
145 text). This behavior is also not related to the actinide boost phenomenon that is observed in some
146 r-process-enhanced stars (supplementary text, Fig. S3). It cautions against using [Ba/Eu] ratios to
147 discern the presence of small amounts of s-process material in stars where the r-process
148 contribution is dominant (supplementary text). The correlations are more closely associated with
149 [Eu/Fe] than $\log \epsilon(\text{Eu})$ as a measure of the strength of the r-process (supplementary text), which

150 suggests a link between r-process production of fissioning transuranic elements and the local
151 environment.

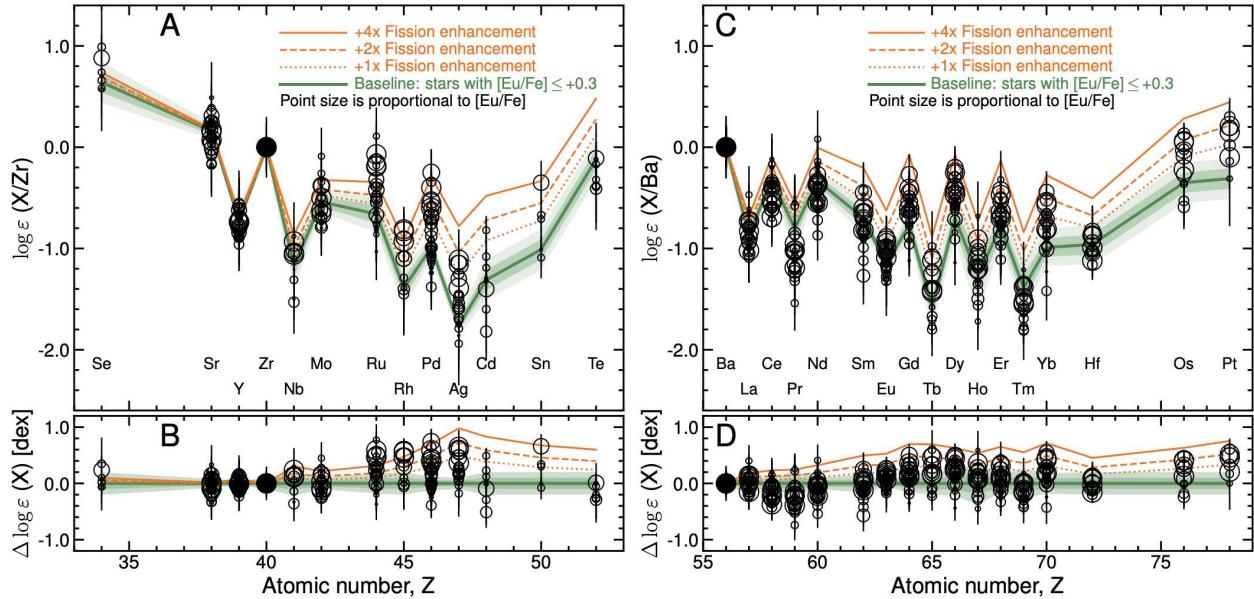
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153 Many of the radioactive, neutron-rich nuclei populated during the r-process are inaccessible to
154 laboratory experiments, so theoretical models are necessary to estimate their properties. These
155 models must be extrapolated from more stable isotopes located near the valley of stability. This
156 extrapolation leads to large uncertainties in key quantities, such as masses, halflives, neutron-
157 capture cross sections, and decay channels, which propagate into uncertainties in the abundance
158 yields (18). Our dataset provides new constraints on these models. For example, many of the
159 fissioning nuclei influencing r-process abundances are predicted to have asymmetric
160 distributions with a lighter peak and a heavier peak (19). The lighter fragments are likely
161 contributing to the Ru, Rh, Pd, and Ag abundances in these stars. Atomic mass numbers $99 \leq A \leq$
162 110 represent the stable isotopes of these elements that can be synthesized in the r-process, via β
163 decays of the neutron-rich nuclei created directly through fission deposition. The heavier
164 fragments likely contribute to the Eu and heavier elements ($A > 150$) in these stars. These mass
165 ranges imply that nuclei with $A > 260$ (i.e., 110 + 150) are produced in the r-process and
166 contribute substantially to the fission fragments. If the r-process synthesizes nuclei as heavy as A
167 ~ 280 , around the predicted nuclear shell closure at 184 neutrons (13, 20), the heavier-mass
168 fragments would include nuclei with $A > 170$ and encompass the lanthanide elements (19, 21,
169 22). This prediction is consistent with our findings. Much larger samples of stars with a wide
170 range of levels of r-process enhancement—and additional observations of Cd, Sn, Te, Hf, Os,
171 and Pt in particular (supplementary text)—are needed to further constrain these model
172 predictions.

173

174 The observational link between elements with $44 \leq Z \leq 47$ and $63 \leq Z \leq 78$ offers the strongest
175 evidence yet for r-process production of nuclei with $A > 260$. Our findings boost the prospects
176 for searches for the γ -rays associated with fission fragments (23) and efforts to assess how
177 fissioning species impact the kilonova light curves associated with neutron star mergers (24–26).
178 They can motivate new theoretical efforts to determine whether superheavy elements are
179 produced in the r-process (27). Astronomical observations of fission fragments can help
180 prioritize the most critical measurements of the properties of short-lived neutron-rich nuclei with
181 the next generation of radioactive-isotope beam facilities (28). This new line of inquiry can also
182 complement and focus studies of the fission debris from the Trinity test (29) and other tests
183 conducted on nuclear explosion debris (30, 31), as well as studies of fission products extracted
184 from the natural reactor located at Oklo, Gabon (32). The fission products found in r-process-
185 enhanced stars may even probe much heavier parent nuclei than ones accessible to these
186 terrestrial sources. Our study highlights the importance of collecting new astronomical data in
187 pursuit of this goal.

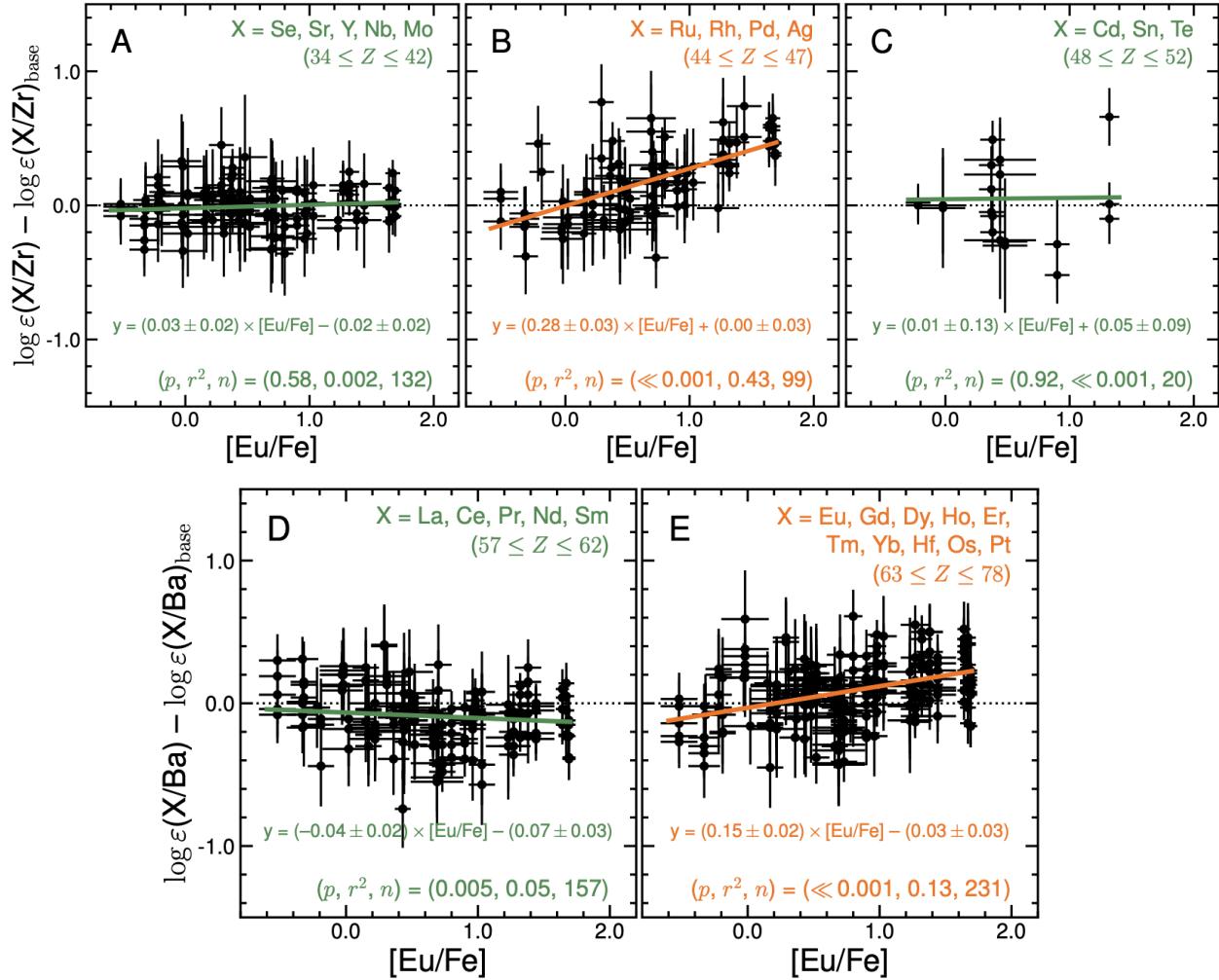
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191 **Fig. 1: Comparison between the observed abundance patterns and fission model**

192 **predictions.** The baseline pattern (green line) is defined as the mean ± 1 or 2 times the standard
 193 error (shaded green bands) of the abundance ratios in the stars with $[\text{Eu}/\text{Fe}] \leq +0.3$. The orange
 194 curves illustrate fission fragments added to this baseline pattern; e.g., $+2\times$ enhancement means 2
 195 parts fission fragments and 1 part baseline pattern. The size of each point is proportional to the
 196 $[\text{Eu}/\text{Fe}]$ ratio in the star, as shown in Fig. 1. Panels **A** and **C** illustrate the $\log \varepsilon$ abundance ratios,
 197 defined as $\log_{10}(N_X/N_{\text{Zr}})$ or $\log_{10}(N_X/N_{\text{Ba}})$, where N_X , N_{Zr} , and N_{Ba} represent the number densities
 198 of elements X (for $X = \text{Se}, \text{Sr}, \dots, \text{Te}$) and Zr or X (for $X = \text{La}, \text{Ce}, \dots, \text{Pt}$) and Ba , respectively.
 199 Panels **B** and **D** illustrate the differences between each star or fission model and the baseline
 200 pattern.

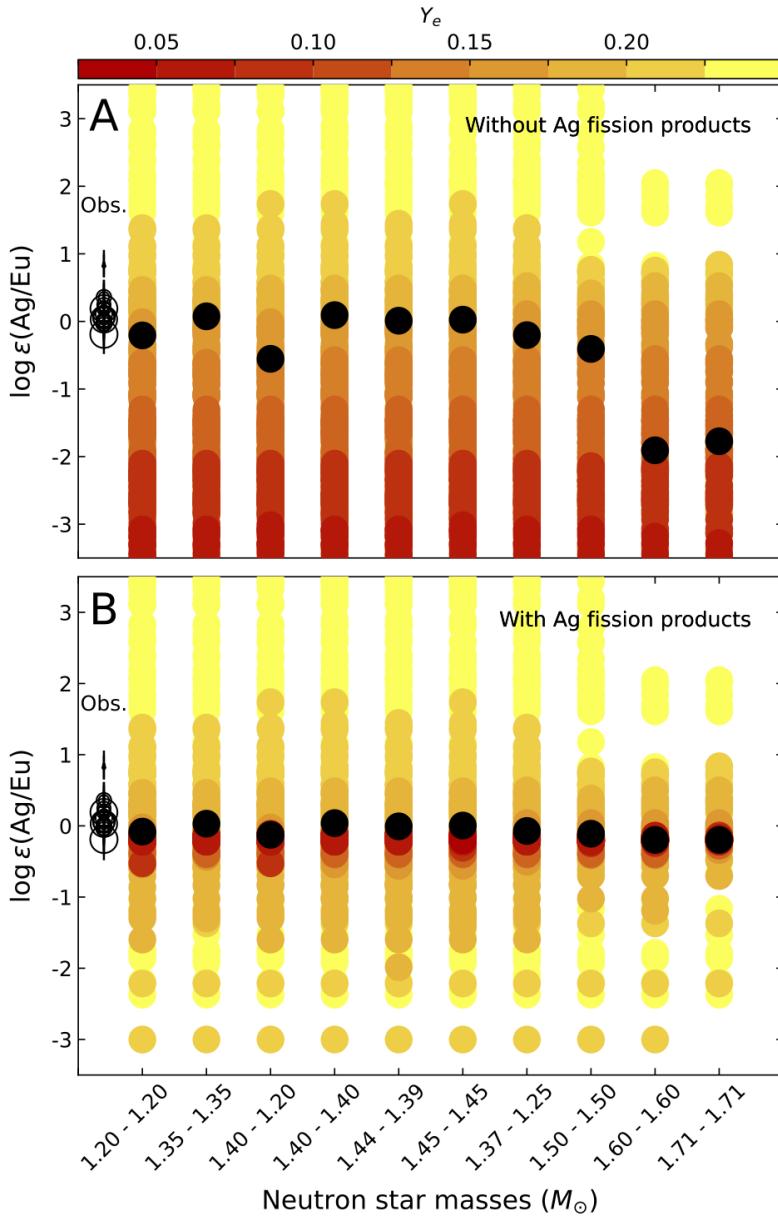
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202 **Fig. 2: Chemical abundance ratios for elements that exhibit no significant correlations with**
 203 **$[Eu/Fe]$ (panels A, C, and D) and those that do (panels B and E).** The panels are ordered by
 204 increasing atomic number, Z . Each point represents one ratio in one star, and the errorbars
 205 represent $\pm 1\sigma$ uncertainties. The $\log \varepsilon(X/Zr)$ and $\log \varepsilon(X/Ba)$ ratios are normalized to the
 206 baseline patterns, $\log \varepsilon(X/Zr)_{\text{base}}$ and $\log \varepsilon(X/Ba)_{\text{base}}$. The abundances of the lighter r-process
 207 elements ($34 \leq Z \leq 52$) are only partially correlated with the abundances of the heavier r-process
 208 elements ($56 \leq Z \leq 78$), so each group is normalized separately to Zr and Ba, respectively. Solid
 209 lines mark weighted least-squares linear fits, and the equations for these fits are printed in each
 210 panel. Dotted lines mark differences of zero. Flat trends (green lines) indicate that an element is
 211 coproduced with Zr or Ba, and significant correlations (orange lines) indicate that an element is

212 coproduced with Eu. The three values printed in the lower right corner of each panel list the *p*-
213 value for the Pearson correlation coefficient, the r^2 coefficient of determination, and the number
214 of stars, *n* (8).

215
216



217 **Fig. 3: The $\log \epsilon(\text{Ag}/\text{Eu})$ abundance ratio predicted when fission deposition contributes to**
 218 **Ag (panel B) and when it does not (panel A).** Results from hydrodynamic simulations (8) for
 219 neutron star mergers with various progenitor masses are shown as colored dots for individual
 220 ejecta and black dots for the mass-weighted abundance ratio. Individual ejecta are colored based
 221 on their neutron richness, Y_e , with lower Y_e ejecta able to reach fissioning nuclei. When fission
 222 deposition contributes to Ag, the predicted mass-weighted $\log \epsilon(\text{Ag}/\text{Eu})$ ratios are independent of

223 the progenitor masses and match the observed ratios (black circles, whose size is proportional to
224 the [Eu/Fe] ratio in each star).

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226 **References and notes**

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981

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983 available in the main text or supplementary materials. The nucleosynthesis yields are
984 available at this Zenodo link: <https://zenodo.org/record/7127232> .

985

986

987 **Supplementary Materials**

988

989 Materials and Methods

990 Supplementary Text

991 Figs. S1 to S3

992 Tables S1 to S6

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Supplementary Materials for

Title: Observational signatures of transuranic fission fragments in stars

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This PDF file includes:

- Materials and Methods
- Supplementary Text
- Figs. S1 to S3
- Tables S1 to S6

Other Supplementary Materials for this manuscript include the following:

N/A

27 **Materials and Methods**

28

29 Stellar sample

30

31 We assemble our sample from the stellar abundance literature of the last few decades. Our
32 sample is drawn from stars listed in the JINAbase stellar abundance database (33), supplemented
33 with our own knowledge of studies not included in JINAbase. We require that Zr, Ba, and Eu
34 abundances are reported (i.e., detected) for each star. We require that $[\text{Ba/Eu}] < -0.3$ for each
35 star, which indicates that r-process material dominates the heavy-element abundance pattern
36 (34). This cut effectively excludes stars with contributions from other processes, such as the s-
37 process or i-process, and, as discussed in detail below, it does not interfere with one of the
38 deviations from r-process universality (subtle variations in the [Ba/Eu] ratio). Finally, we require
39 that an abundance of at least one of the elements Se, Pd, or Te is reported, because we are
40 interested in studying the abundances of elements at and between the first and second r-process
41 peaks. Our sample is complete within these requirements to the best of our knowledge.

42

43 Our sample is comprised of 42 stars spanning ≈ 2.6 dex in metallicity ($-3.57 \leq [\text{Fe/H}] \leq -0.99$)
44 and ≈ 2.2 dex in r-process enhancement ($-0.52 \leq [\text{Eu/Fe}] \leq +1.69$). Table S1 lists the name of
45 each star, its metallicity ($[\text{Fe/H}]$), its r-process enhancement ($[\text{Eu/Fe}]$), abundances for the
46 elements Se through Te, and the literature references from which we have adopted the
47 abundances. Table S2 lists the uncertainties in these abundances. Table S3 lists the abundances
48 of the heavier r-process elements, Ba through Th, in these stars, and Table S4 lists the
49 uncertainties in these abundances.

50

51 We homogenize the stellar abundances from the literature as follows. Stellar abundances depend
52 directly on the atomic transition probability of each line, expressed as the $\log(gf)$ value, which is
53 the log of the degeneracy of the lower level times the oscillator strength. Different stellar
54 abundance studies may adopt different sets of $\log(gf)$ values for the same set of transitions,
55 which artificially increases the statistical scatter in the set of combined abundances. Furthermore,
56 laboratory measurements of the quantities necessary to calculate $\log(gf)$ values have been
57 improved in recent years for transitions of some of the lighter r-process elements, which can
58 improve the accuracy and precision of the stellar abundances. We therefore translate the
59 abundances of Se through Te to a uniform $\log(gf)$ scale, with references given in (46). These
60 changes are always small, ≤ 0.14 dex, and typically ≤ 0.03 dex, so the impact on the literature
61 abundances is, at most, minor. The updated $\log \varepsilon(X)$ abundances are presented in Table S1. The
62 heavier r-process elements are generally drawn from studies that are based on modern laboratory
63 measurements (71), so no corrections are made for the abundances of these elements.

64

65 Some heavy elements (Br, Kr, Rb, In, Sb, I, Xe, Cs, Lu, Ta, W, Re, Ir, Au, Hg, Tl, Pb, Bi, and U)
66 are omitted from our analysis because there are insufficient observations of them. Tc, Pm, Po,
67 At, Rn, Fr, Ra, Ac, and Pa are omitted because they have no stable or long-lived isotopes. Th is
68 discussed in a separate section below.

69

70

71 Abundance correlations and non-correlations in our sample

72

73 Figures S1 and S2 illustrate the relationships among each of the lighter r-process elements ($34 \leq$
74 $Z \leq 52$) and the heavier r-process elements ($56 \leq Z \leq 78$) with the [Eu/Fe] ratios. Weighted linear
75 least-squares fits are shown.

76

77 Figure S1 demonstrates that the correlations between the $\log \varepsilon(X/Zr)$ ratios and [Eu/Fe] are not
78 significant ($p \geq 0.08$, whereas $p < 0.05$ would indicate significance) for $X = \text{Se, Sr, Y, Nb, Mo,$
79 Cd, Sn, and Te. The relationships are significant ($p \leq 0.002$) for $X = \text{Ru, Rh, Pd, and Ag}$. Figure
80 S2 demonstrates that no significant correlations are found among the $\log \varepsilon(X/Ba)$ ratios and
81 [Eu/Fe] for $X = \text{La, Nd, Sm, Tb, Ho, Tm, Hf, Os, and Pt}$. Significant ($p \leq 0.03$) correlations are
82 found for $X = \text{Ce, Pr, Eu, Gd, Dy, Er, and Yb}$. Hints of similar positive correlations are
83 suggested by the data for heavier r-process elements with limited numbers of measurements (Tb,
84 Ho, Tm, Hf, Os, Pt; $n \leq 21$ stars). We flag these elements as being ripe for future observational
85 studies to confirm or refute the positive trends hinted at by current data. The significant
86 correlations for Ce and Pr ($Z = 58$ and 59) are in the opposite sense as the correlations we
87 associate with signatures of fission fragments. In other words, their $\log \varepsilon(X/Ba)$ versus [Eu/Fe]
88 slopes are negative, and they are consistent with zero within 2 standard deviations (σ). The cause
89 of this behavior is unclear at present, and we also flag it as being ripe for future theoretical
90 studies to investigate.

91

92 This behavior justifies our decision to collect these elements into five groups, as shown in Fig. 2:
93 $34 \leq Z \leq 42$ (Panel A; no significant correlations with [Eu/Fe]), $44 \leq Z \leq 47$ (Panel B; significant
94 positive correlations with [Eu/Fe]), $48 \leq Z \leq 52$ (Panel C; no significant correlations with

95 [Eu/Fe]), $56 \leq Z \leq 62$ (Panel D; no significant positive correlations with [Eu/Fe]), and $63 \leq Z \leq$
96 78 (Panel E; significant positive correlations, or hints thereof, with [Eu/Fe]).

97
98 We next assess the statistical significance of the correlations within each of these five element
99 groups. For the groups shown in Panels A, C, and D of Fig. 2, the slope of the weighted linear
100 least-squares fit is consistent with zero within 2σ for all three sets of elements. The r^2 (coefficient
101 of determination) values for the elements shown in Panels A, C, and D indicate that only 0.2%,
102 $\ll 0.1\%$, and 5%, respectively of the variations in the abundances can be associated with the
103 correlation with [Eu/Fe], rather than random uncertainties in the abundances. Furthermore, the p -
104 value for the Pearson correlation coefficient for each sample cannot reject the null hypothesis of
105 zero slope at even modest significance ($p \geq 0.05$) for two of the three sets of elements; the p -
106 value for the third set (Ba through Sm; Panel D) is 0.005 for 157 measurements, although the
107 slope is negative and consistent with zero (-0.04 ± 0.02) within 2σ .

108
109 The groups shown in Panels B and E of Fig. 2 exhibit different behavior. Their slopes differ from
110 zero by more than 7σ . The r^2 values for the elements shown in Panels B and E indicate that 43%
111 and 13%, respectively, of the variations in the abundances can be associated with the correlation
112 with [Eu/Fe]. Their p -values are also highly significant (Ru through Ag: $p \ll 0.001$ for $n = 97$
113 individual measurements; Eu through Pt: $p \ll 0.001$ for $n = 231$ individual measurements).

114 These metrics collectively suggest that the relationships shown in Panels B and E of Fig. 2 are
115 statistically significant.

116

117

118 Nucleosynthesis calculations: hydrodynamics, nuclear data, and fission model

119

120 The nucleosynthesis calculations presented in this work make use of the hydrodynamic
121 simulations of binary neutron star merger dynamical ejecta from (72), which considers a number
122 of variations including the neutron star equation of state, differing treatments of neutrino
123 transport, as well as distinct sets of progenitor masses. Here we use the simulation results
124 reported for the case of the LS220 equation of state with a neutrino leakage scheme, since these
125 are the conditions for which (72) considered a broad set of progenitor mass variations. Our
126 adoption of these models does not necessarily imply that neutron star mergers are the only viable
127 site of r-process nucleosynthesis. Rather, this scenario provides a reasonable set of conditions in
128 which to explore the behavior of fissioning transuranic nuclei in an r-process environment.

129

130 We demonstrate the stabilizing effect fission deposition has on predicted abundance ratios in Fig.
131 3. We consider two distinct fission yield sets for neutron-rich nuclei. In the first one, we assume
132 a fissioning species always splits in half, concentrating products near $A \sim 130$ and producing no
133 lighter elements, such as Ag. In the second one, we adopt a proper theoretical model for the
134 fission yields (73), which makes use of the properties predicted by the finite range liquid drop
135 model (FRLDM) for very heavy unstable species. This model produces a significant amount of
136 fission products with $A < 130$.

137

138 When exploring variations in the fission yields, we keep all other input nuclear data the same and
139 use the datasets described in (16). These datasets include mass predictions from FRDM2012,

140 with neutron capture and neutron-induced fission rates, and β -decay and β -delayed fission rates,
141 determined from the LANL CoH and BeoH models, respectively. We also adopt the FRLDM
142 fission barriers. Although significant uncertainties remain in the nuclear data for the neutron-rich
143 nuclei populated during our nucleosynthesis calculations, and such unknowns can impact the
144 exact abundance ratios predicted (74), our result regarding fission deposition and coproduction
145 does not depend on our choice of adopted nuclear data. Rather, the critical requirements
146 highlighted by these calculations are (a) that the r-process reaches nuclei that fission during the
147 times when abundances are being set, and (b) that these fission products include both lighter
148 elements, such as Ag, and heavier elements, such as Eu, thereby linking these elements through
149 their simultaneous production.

150

151

152 Modeling the abundance patterns with the inclusion of fission fragments

153

154 Our analysis of the abundance pattern that includes enhancements from fission fragments is
155 constructed as follows. Our empirical template for the baseline abundance pattern is calculated
156 using the mean abundance ratios found in the 13 stars with $[\text{Eu}/\text{Fe}] \leq +0.3$. We assume that these
157 stars contain minimal contributions from the process responsible for producing the abundance
158 excesses observed in Figs. 1 and 2. This assumption is justified because the star-to-star
159 dispersion of the abundance ratios in these stars (median σ of 0.19 dex) is comparable to the
160 typical observational uncertainties (median uncertainty of 0.22 dex). In other words, the
161 abundance ratio dispersion in these 13 stars can be fully accounted for by the observational
162 uncertainties.

163

164 The values of this baseline pattern, expressed as $\log \varepsilon(X/Zr)_{\text{base}}$ and $\log \varepsilon(X/Ba)_{\text{base}}$, are listed in
165 Tables S5 and S6. The pattern is shown by the green line in Fig. 1. We assume that this baseline
166 r-process pattern is present in all stars in the sample. We add enhancements of transuranic fission
167 fragments, adopted from the models described in the previous section. We scale the level of
168 fission fragment enhancement, F_{enh} , to the $\log \varepsilon(Pd/Zr)$ ratio for the lighter r-process elements
169 and the $\log \varepsilon(Gd/Ba)$ ratio for the heavier r-process elements. For example, $F_{\text{enh}} = 2$ indicates 1
170 part baseline and 2 parts fission fragments. Mathematically, the abundance ratio for element X
171 relative to Zr can be expressed as

172

173 $\log \varepsilon(X/Zr)_{\text{total}} = \log(F_{\text{enh}} \times 10^{\log \varepsilon(X/Zr)_{\text{fiss}}} + 10^{\log \varepsilon(X/Zr)_{\text{base}}}),$

174

175 where

176

177 $\log \varepsilon(X/Zr)_{\text{fiss}} = \log \varepsilon(X)_{\text{fiss}} - \log \varepsilon(Pd)_{\text{fiss}} + \log \varepsilon(Pd/Zr)_{\text{base}}.$

178

179 The values of $\log \varepsilon(X)_{\text{fiss}}$ and $\log \varepsilon(Pd)_{\text{fiss}}$ are adopted from the fission model with equal-mass
180 $1.40 M_{\odot}$ merging neutron stars, although the exact model selected has little influence on the
181 result, as shown in Fig. 3. All logarithms are base 10. We use an analogous set of equations to
182 describe the enhancements relative to the baseline abundance pattern for the heavier r-process
183 elements, replacing Zr with Ba and Pd with Gd. The enhancements described by these equations
184 are illustrated by the orange curves in Fig. 1.

185

186 We estimate the relative contributions from fission fragment deposition under the assumption
187 that stars with $[\text{Eu}/\text{Fe}] \leq +0.3$ contain minimal contributions from fission. According to the
188 weighted least-squares linear fit shown in Fig. 2, the Ru through Ag abundances increase by a
189 factor of ≈ 2 (0.30 dex) in stars with $[\text{Eu}/\text{Fe}] \approx +1.1$. Thus, roughly half of the Ru through Ag in
190 stars with $[\text{Eu}/\text{Fe}] \approx +1.1$ originated as fission fragments. These elements are enhanced by an
191 average factor of ≈ 3 (0.48 dex) in the stars with the highest r-process enhancement, revealing
192 that fission may be responsible for up to $\approx 3/4$, or $\approx 75\%$, of the Ru through Ag in these stars.
193 Similarly, the Eu through Pt abundances are enhanced by an average factor of ≈ 1.7 (0.22 dex) in
194 the stars with the highest r-process enhancement, revealing that fission may be responsible for up
195 to $\approx 60\%$ of the Eu and heavier lanthanide elements in these stars.

196

197

198

199 **Supplementary Text**

200

201 Previous investigations of deviations from r-process abundance universality

202

203 R-process universality has been generally understood to apply to heavy r-process elements,
204 including those with $56 \leq Z \leq 72$ (37) and $76 \leq Z \leq 78$ (34). Two important deviations from r-
205 process universality have been identified in observations. The larger effect of the two, whose
206 variation may span more than 1.5 dex from one metal-poor star to another, is characterized by
207 differing abundance levels between the lighter ($34 \leq Z \leq 52$) and heavier ($Z \geq 56$) r-process

208 elements (49, 57, 62, 75–79). The abundances of the two element groups are not fully
209 independent of one another. The abundance pattern within each of the two groups is generally
210 consistent, once the overall abundance level has been normalized (37, 55, 56, 68), with the
211 exception of the excesses we have identified. The second effect, whose variation may span up to
212 0.7 dex, is characterized by differing abundance levels between the actinides (Th and U; $Z = 90$
213 and 92) and the heaviest stable r-process elements (20, 35, 80–83). This effect, known as the
214 actinide boost when the actinides are enhanced, is discussed in more detail in a separate section
215 below. These deviations from universality may reflect a diversity of conditions within a given r-
216 process site (11, 83), multiple sites within events (84–87), multiple events (10), or some
217 combination of these scenarios.

218

219

220 Other evidence for cosmic production of transuranic nuclei

221

222 Studies have found evidence for the cosmic production of the transuranic nuclei up to at least $A =$
223 247. The decay products of ^{244}Pu and ^{247}Cm are found in meteorites (88, 89). ^{244}Pu has also been
224 detected directly in seafloor sediments (90, 91). Transuranic nuclei have not been observed in
225 stars, and their halflives are short (≤ 80 Myr) relative to the ages of stars where they might
226 otherwise be observable ($\gg 1$ Gyr).

227

228

229 Alternative explanations

230

231 The correlations between the $\log \epsilon(\text{Ru/Zr})$, $\log \epsilon(\text{Rh/Zr})$, $\log \epsilon(\text{Pd/Zr})$, and $\log \epsilon(\text{Ag/Zr})$ ratios
232 and $[\text{Eu/Fe}]$ indicates that the source of the elevated Ru, Rh, Pd, and Ag abundances is related to
233 the source of the heavy elements, including Eu. In principle, that source could be a different
234 nucleosynthesis process operating in a core-collapse supernova, such as the weak r-process,
235 weak s-process, or the intermediate neutron-capture process (i-process). We can exclude any
236 process that operates only in low- or intermediate-mass stars that pass through the asymptotic
237 giant branch (AGB) phase of evolution, such as the main s-process (92). Therefore, we compare
238 the observed abundance behavior with predictions for the weak r-process, the weak s-process,
239 and the i-process.

240

241 The weak r-process is associated with core-collapse supernovae, where the matter is only slightly
242 neutron rich. The weak r-process is generally not predicted to produce substantial amounts of
243 elements as heavy as Eu (93–95). Therefore, the weak r-process is not a viable explanation for
244 the observed abundance behavior.

245

246 The weak s-process occurs in massive, rapidly rotating stars. Models of the weak s-process in
247 low-metallicity environments do not predict enhancement among any of the $\log \epsilon(\text{Ru/Eu})$, \log
248 $\epsilon(\text{Rh/Eu})$, $\log \epsilon(\text{Pd/Eu})$, or $\log \epsilon(\text{Ag/Eu})$ ratios without a similar increase in the $\log \epsilon(\text{Ba/Eu})$ and
249 $\log \epsilon(\text{Pb/Eu})$ ratios. These models predict enhancement of the $\log \epsilon(\text{Ba/Eu})$ and $\log \epsilon(\text{Pb/Eu})$
250 ratios by several orders of magnitude (96–98). These values exceed the observed ratios, \log
251 $\epsilon(\text{Ba/Eu}) = 0.99$ ($\sigma = 0.07$ dex) in the nine stars with $[\text{Eu/Fe}] > +1.0$ in our sample, and \log
252 $\epsilon(\text{Pb/Eu}) \approx 0.8$ ($\sigma \approx 0.2$ dex) in other r-process-enhanced stars (79). The weak s-process is also
253 not a viable explanation for the observed abundance behavior.

254

255 The i-process has been associated with several sites, including super-AGB stars, post-AGB stars,
256 He-core and He-shell flashes in low-mass stars, and rapidly accreting white dwarfs. Models of
257 the i-process can produce both Pd and Eu, as well as other heavy elements (99–101). In order to
258 produce $\log \epsilon(\text{Pd/Eu})$ ratios similar to those observed in the nine stars with $[\text{Eu/Fe}] > +1.0$, \log
259 $\epsilon(\text{Pd/Eu}) = 0.71$ ($\sigma = 0.12$ dex), these models predict $\log \epsilon(\text{Ba/Eu})$ ratios $\gtrsim 2$, which is much
260 higher than observed in these nine stars, $\log \epsilon(\text{Ba/Eu}) = 0.99$ ($\sigma = 0.07$ dex). Furthermore, the i-
261 process would need to dominate the production of Eu in the stars with the highest levels of
262 $[\text{Eu/Fe}]$, which is unlikely given that the abundance ratios among Eu, the lanthanide elements,
263 and all heavier elements are consistent with the Solar System r-process ratios and models for r-
264 process nucleosynthesis (102). Therefore, the i-process is not a viable explanation for the
265 observed abundance behavior.

266

267 We conclude that no known nucleosynthesis process, except for the deposition of transuranic
268 fission fragments in the r-process, can explain the observed abundance behaviors.

269

270

271 Using [Ba/Eu] ratios as a metric to assess the relative contributions from the r- and s-process

272

273 A high percentage (>94% or so) of the Eu in the Solar System originated via r-process
274 nucleosynthesis. In contrast, a high percentage of the Ba in the Solar System (> 85% or so)
275 originated via s-process nucleosynthesis. This situation naturally invites the use of the [Ba/Eu]
276 ratio (or, analogously, the $\log \epsilon(\text{Ba/Eu})$ ratio) as a metric to assess the relative contributions of

277 the r- and s-process to different stars. The exact values that define "pure" r- and s-process
278 [Ba/Eu] ratios depend somewhat on the adopted model and method, but they typically differ by >
279 2 dex (103–108). For example, (108) estimates that the [Ba/Eu] ratios for "pure" r-process and
280 "pure" s-process material are ≈ -0.9 and $\approx +1.3$, respectively.

281
282 Our findings imply that the [Ba/Eu] ratio exhibits a small amount of cosmic dispersion due to the
283 r-process itself. The eight stars in our sample with $[\text{Eu}/\text{Fe}] < 0.0$ exhibit $[\text{Ba}/\text{Eu}] = -0.53 \pm 0.05$
284 ($\log \epsilon(\text{Ba}/\text{Eu}) = 1.13 \pm 0.05$). The nine stars with $[\text{Eu}/\text{Fe}] > +1.0$ exhibit $[\text{Ba}/\text{Eu}] = -0.67 \pm 0.02$
285 ($\log \epsilon(\text{Ba}/\text{Eu}) = 0.99 \pm 0.02$). These differences are small but significant at about the 2.5σ level.

286
287 Fortunately, the [Ba/Eu] ratio remains an acceptable general diagnostic of the r- and s-process
288 contributions. The differences between the low- and high-[Eu/Fe] samples are much smaller than
289 the > 2 dex differences in the "pure" r- and s-process abundance ratios. Caution is warranted,
290 however, when attempting to interpret slight (< 0.2 dex) enhancements in the [Ba/Eu] ratio as
291 necessarily implying the presence of small amounts of s-process contamination to an otherwise
292 dominant r-process pattern. Similar cautions apply to other element ratios, such as [La/Eu] (109)
293 or [Ce/Eu], where small but measurable enhancements could also signal deficiencies in Eu that
294 result from less fission fragment deposition.

295
296
297 A metric to assess the strength of the r-process
298

299 Reference (16) compared the predictions of their fission yields with a limited set of observational
300 data, 13 r-process-enhanced stars drawn from JINABase (33). That study examined the
301 relationships between the [Ru/Eu], [Pd/Eu], and [Ag/Eu] ratios and the $\log \varepsilon(\text{Eu})$ ratio, which
302 they adopted as a measure of the level of r-process enrichment in a given star. That study noted
303 no correlations between the [Ru/Eu], [Pd/Eu], or [Ag/Eu] ratios and $\log \varepsilon(\text{Eu})$, which was
304 consistent with the behavior expected if the Ru, Pd, and Ag originate in part as fission fragments.

305 We repeat a similar exercise for the stars in our sample, and we find significant correlations
306 between $\log \varepsilon(\text{Eu})$ and each of $\log \varepsilon(\text{Ru/Zr})$ ($p = 0.03$, 29 stars) and $\log \varepsilon(\text{Ag/Zr})$ ($p = 0.02$, 21
307 stars). These correlations, while significant, are weaker than the correlations with [Eu/Fe].

308

309 We find that [Eu/Fe], rather than $\log \varepsilon(\text{Eu})$ (or [Eu/H]), represents a more predictive metric of
310 the strength of the r-process. It is unclear at present why this is the case. The $\log \varepsilon(\text{Eu})$ ratios
311 represent the amount of Eu and r-process elements present in the gas from which these stars
312 formed. In contrast, the [Eu/Fe] ratios represent the amount of Eu and r-process elements present
313 relative to the Fe abundance of the gas, which suggests a connection to the environment where
314 the r-process occurred (110–112). Future theoretical attempts to explore the nature and cause of
315 this behavior will be an important step to establish the relationship between the production of
316 transuranic elements and the site(s) of the r-process.

317

318

319 No relation to the actinide boost phenomenon

320

321 The actinide boost phenomenon is not correlated with the production of transuranic fission
322 fragments in our data, as shown in Fig. S3, panel A. We adopt $\log \varepsilon(\text{Th/Eu})$ as a measure of the
323 range of actinide production relative to the lanthanide elements. We adopt $\log \varepsilon(\text{Pd/Zr})$ as a
324 measure of the fission fragment yields relative to the lighter r-process elements. For the 15 stars
325 in the sample with reported Pd and Th abundances, the *p*-value for the Pearson correlation
326 coefficient between the $\log \varepsilon(\text{Th/Eu})$ and $\log \varepsilon(\text{Pd/Zr})$ ratios is 0.90, indicating a non-significant
327 correlation. Furthermore, the variation in the observed actinide abundances relative to the
328 lanthanide abundances is much smaller than the variation observed among the fission fragments.
329 The ranges of the $\log \varepsilon(\text{Th/Eu})$ and $\log \varepsilon(\text{Pd/Zr})$ ratios are 0.45 dex and 1.13 dex, respectively,
330 which correspond to factors of 2.8 and 13.5. The fission fragment yields vary by ≈ 5 times more
331 than the actinide variations.

332

333 As shown in Fig. S3, panel B, the relationship between the fission fragments and the actinides is
334 also not correlated with the strength of the r-process. We adopt $\log \varepsilon(\text{Pd/Th})$ as a measure of the
335 amount of fission fragments relative to the actinides. We adopt $[\text{Eu/Fe}]$ as a measure of the
336 strength of the r-process. The $\log \varepsilon(\text{Pd/Th})$ ratio is not correlated with the $[\text{Eu/Fe}]$ ratio. The *p*-
337 value for the Pearson correlation correlation coefficient between the $\log \varepsilon(\text{Pd/Th})$ and $[\text{Eu/Fe}]$
338 ratios is 0.37 for the 15 stars in our sample, indicating a non-significant correlation.

339

340 We conclude that there is no correlation between the fission fragment yields and the actinide
341 yields in these stars.

342

343

344 Limitations

345

346 The baseline abundance pattern in our analysis (Fig. 1, Tables S5 and S6) is calculated
347 empirically, rather than self-consistently with the fission component. Our analysis preserves the
348 observed abundance ratios among Se, Sr, Y, Zr, Nb, Mo, Rh, and Pd, as shown in Fig. 1. It
349 underpredicts the abundances for Ru and overpredicts the abundances of Ag, Cd, Sn, and Te,
350 revealing that there is room for improvement in the models. For example, there could be a more
351 narrow range of nuclei undergoing fission during the r-process, which could therefore narrow the
352 range of daughter products produced. Future model explorations should aim to reproduce the full
353 range of observed abundance behavior self-consistently, and our results signal that astronomical
354 observations can be used to confront those models with observations.

355

356 Our conclusion that some elements originate in part via transuranic fission fragment deposition
357 in the r-process may not be applicable to populations that are more metal-rich or contaminated
358 by material from other processes, such as the s-process. A comparison with one particular
359 previous study demonstrates this limitation. Ref. (57) examined the abundances of Sr, Y, Zr, Pd,
360 Ag, Ba, and Eu in 71 metal-poor stars to trace the contributions from different nucleosynthesis
361 processes to each of these elements. Taken at face value, that study reached the opposite
362 conclusion from ours, in that it found that Pd and Ag correlated more strongly with Zr than with
363 Eu. A closer examination, however, reveals that the two studies are probing different populations
364 of stars. First, the stars in our sample are much more metal-poor than those in (57); the median
365 [Fe/H] in our sample is -2.62 , while for the sample in (57) it is -1.84 . Secondly, the stars in our
366 sample are more enhanced in r-process elements than the stars in (57); the median [Eu/Fe] in our

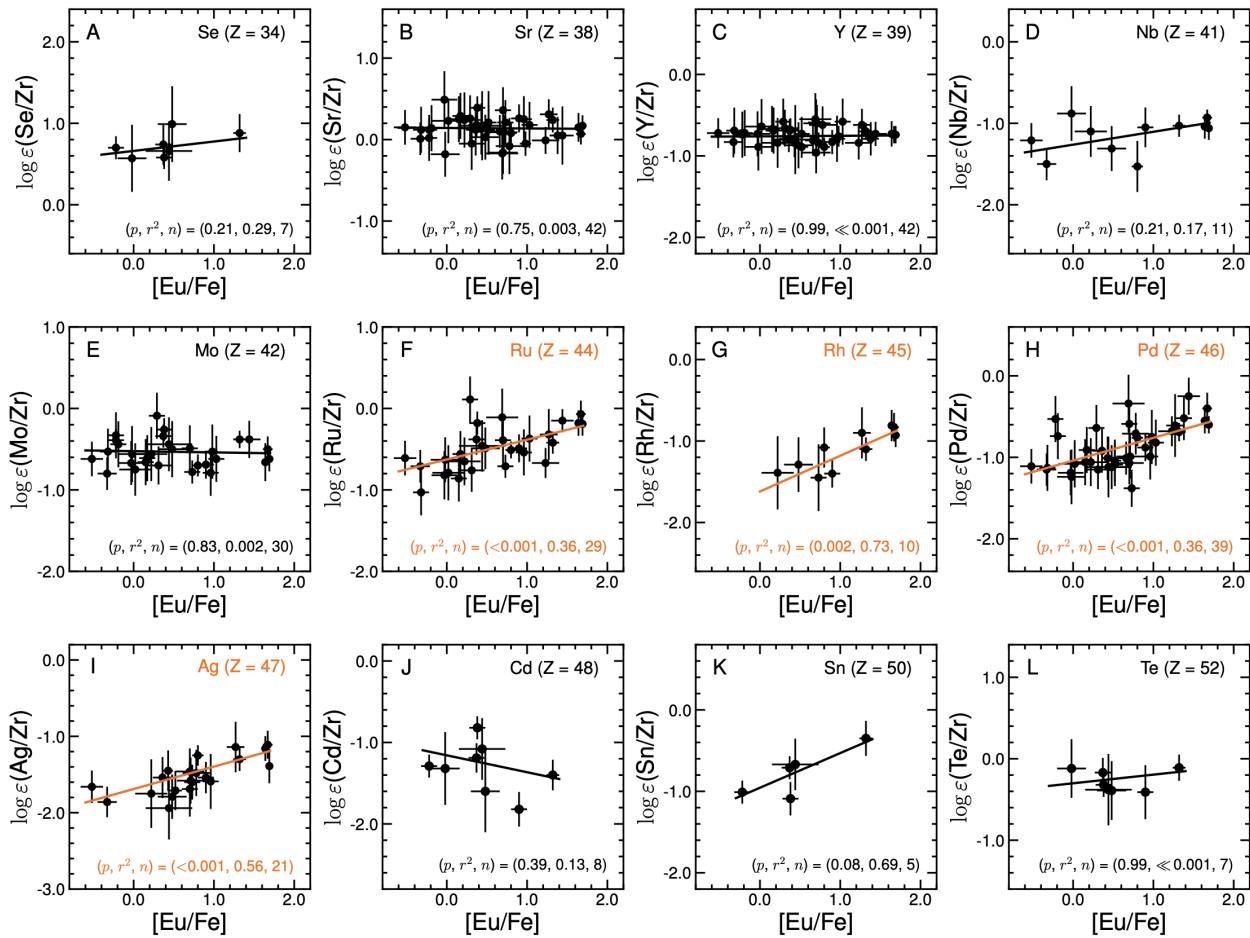
367 sample is +0.52, while for the sample in (57) it is +0.43. Thirdly, the stars in our sample are less
368 contaminated by s-process material than the stars in (57). The median [Ba/Eu] in our sample is
369 –0.55, while for the sample in (57) it is –0.21. The sample examined by (57) was assembled to
370 include stars that would reflect a combination of nucleosynthesis processes, such as charged-
371 particle nucleosynthesis, the weak r-process, the s-process, and the r-process. In contrast, our
372 sample is assembled to include stars where the r-process is the dominant source of the heavy
373 elements.

374

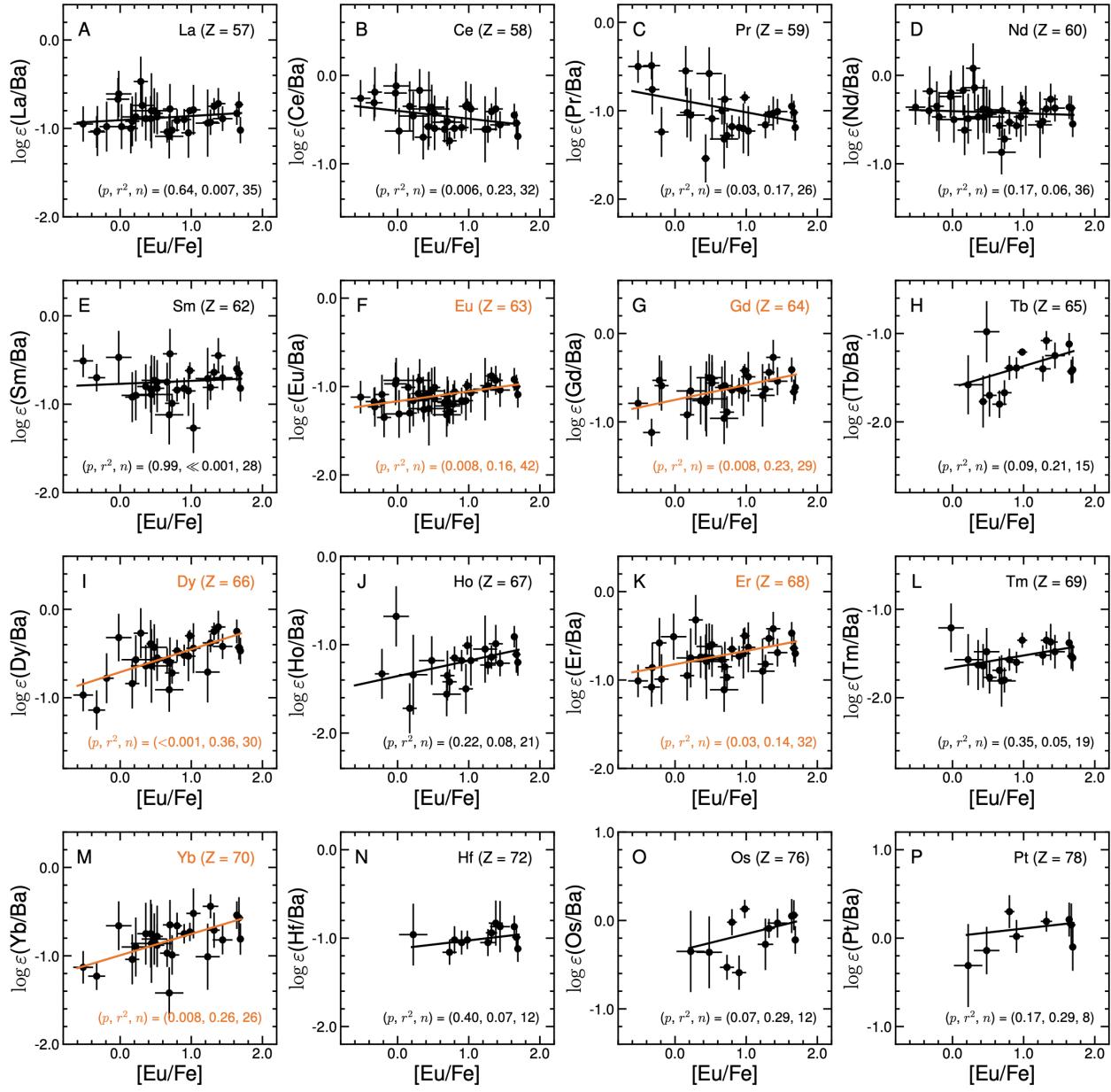
375 Observations of r-process elements in stars are limited to elemental abundances. Individual
376 isotopic abundances are generally inaccessible, with only a few exceptions (113, 114). Our
377 finding that Eu and Ba are slightly decoupled when transuranic fission fragment deposition
378 occurs suggests that measurements of the Ba isotopes in very metal-poor stars (115–118) may
379 need to be reinterpreted.

380

381 Not all elements can be detected in all stars. For example, there are very few observations of Cd,
382 Sn, Te, Hf, Os, and Pt, which limits our ability to draw reliable conclusions about their origin.
383 This limitation is a consequence of several factors, most notably the conditions found in stellar
384 atmospheres and the energy spacing of the electronic configurations. These factors often restrict
385 detection of these elements to wavelengths in the ultraviolet region of the spectrum observable
386 only from space (46). Surveys to identify and characterize many more r-process-enhanced stars
387 by the R-Process Alliance (119–122) and others (123), both today and with the next generation
388 of optical and ultraviolet spectrographs (124–127), will be essential to realize the full potential of
389 these observational constraints.



390 **Fig. S1: Abundance ratios among the lighter elements for the 42 stars in the sample.** Each
 391 point represents one star, and the errorbars represent 1σ uncertainties. Each panel illustrates the
 392 abundance ratios for a different element X (where X = Se, Sr, ..., Te) arranged by increasing
 393 atomic number, Z. Lines mark weighted least-squares linear fits. Non-significant trends (black
 394 lines) indicate that an element is coproduced with Zr, and significant positive correlations
 395 (orange lines) indicate that an element is coproduced with Eu. The three values printed in the
 396 lower right corner of each panel list the p -value for the Pearson correlation coefficient, the r^2
 397 coefficient of determination, and the number of stars, n .

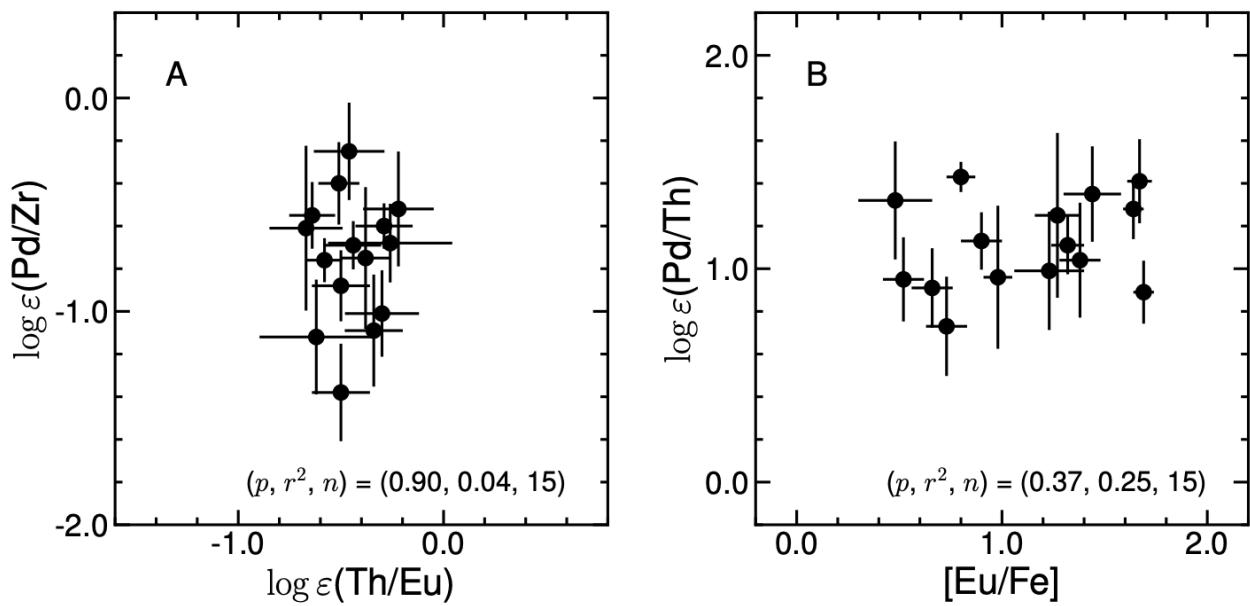


398

399 **Figure S2: Abundance ratios among the heavier elements for the 42 stars in the sample.**

400 Ratios among the lanthanides (La, Ce, Pr, Nd, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb), and
 401 elements near and at the third r-process peak (Hf, Os, Pt) are shown. Each point represents one
 402 star, and the errorbars represent 1σ uncertainties. The panels are arranged by increasing atomic
 403 number, Z . Lines mark weighted least-squares linear fits. Non-significant trends are shown with

404 black lines, and significant trends are shown with orange lines. The three values printed in the
405 lower right corner of each panel list the p -value for the Pearson correlation coefficient, the r^2
406 coefficient of determination, and the number of stars, n .



407 **Figure S3: The relationships between the fission fragment yields (Pd) and the actinides**

408 **(Th).** In panel **A**, the $\log \varepsilon(\text{Pd}/\text{Zr})$ ratio represents the relationship between the fission fragment
 409 yields and lighter r-process elements, and the $\log \varepsilon(\text{Th}/\text{Eu})$ ratios represent the relationship
 410 between the actinides and the lanthanides. In panel **B**, the $\log \varepsilon(\text{Pd}/\text{Th})$ ratios represent the
 411 relationship between the fission fragment yields and the actinides, and the $[\text{Eu}/\text{Fe}]$ ratios
 412 represent the strength of the r-process. All axes span 2.4 dex. The three values printed in the
 413 lower right corner of each panel list the *p*-value for the Pearson correlation coefficient, the r^2
 414 coefficient of determination, and the number of stars, *n*. Neither correlation is significant.
 415

Star	[Fe/H]	[Eu/Fe]	$\log \varepsilon$ (Se)	$\log \varepsilon$ (Sr)	$\log \varepsilon$ (Y)	$\log \varepsilon$ (Zr)	$\log \varepsilon$ (Nb)	$\log \varepsilon$ (Mo)	$\log \varepsilon$ (Ru)	$\log \varepsilon$ (Rh)	$\log \varepsilon$ (Pd)	$\log \varepsilon$ (Ag)	$\log \varepsilon$ (Cd)	$\log \varepsilon$ (Sn)	$\log \varepsilon$ (Te)	Reference
CS 31082-001	-2.90	+1.69		0.68	-0.23	0.51	-0.55	-0.11	0.32	-0.42	-0.09	-0.88				(35–38)
CS 29497-004	-2.84	+1.67		0.72	-0.09	0.65	-0.28	0.15	0.58	-0.18	0.25	-0.46				(39)
CS 22892-052	-3.09	+1.64		0.41	-0.47	0.26	-0.78	-0.40	0.08	-0.55	-0.29	-0.90				(37, 40–42)
J1432-4125	-2.97	+1.44		0.18	-0.60	0.13		-0.25	-0.02		-0.12					(43)
HE 1219-0312	-2.97	+1.38		0.32	-0.52	0.27					-0.25					(44)
HD 222925	-1.46	+1.32	2.62	1.98	1.04	1.74	0.71	1.36	1.32	0.64	1.05	0.44	0.34	1.39	1.63	(45, 46)
J1538-1804	-2.09	+1.27		1.20	0.27	0.89			0.57	-0.01	0.28	-0.25				(47)
CS 31078-018	-2.84	+1.23		0.31	-0.52	0.32			-0.35		-0.36					(48)
CS 22953-003	-2.84	+1.03		0.30	-0.46	0.12		-0.50	-0.25		-0.70					(49, 50)
HE 2327-5642	-2.79	+0.98		-0.01	-0.72	0.04		-0.49			-0.71					(51)
CS 22896-154	-2.69	+0.96		0.74	-0.30	0.49		-0.30	-0.05		-0.50	-1.10				(49, 50)
BD +17°3248	-2.10	+0.90		1.09	0.00	0.83	-0.22	0.14	0.34	-0.57	-0.05	-0.71	-0.99		0.42	(37, 52–55)
HD 221170	-2.20	+0.80		0.81	-0.16	0.73	-0.80	0.03	0.22	-0.35	-0.03	-0.52				(37, 41, 56)
CD -45°3283	-0.99	+0.78		1.73	1.19	1.81					1.10	0.33				(57)
HD 20	-1.60	+0.73		1.47	0.42	1.26		0.48	0.55	-0.19	-0.12	-0.34				(58)
HD 120559	-1.31	+0.71		1.78	1.09	1.69					0.70	0.11				(57)
HD 3567	-1.33	+0.70		1.44	0.64	1.60					0.53	0.14				(57)
BS 17569-049	-2.88	+0.70		0.35	-0.60	-0.01		-0.50	-0.40		-0.60	-1.70				(49, 50)
J1830-4555	-3.57	+0.69		-1.26	-1.64	-1.09			-1.20		-1.43					(59)
HD 115444	-2.85	+0.66		0.05	-0.86	-0.05					-1.06					(60)
HD 186478	-2.60	+0.52		0.49	-0.52	0.37					-0.67	-1.34				(61, 62)
HD 6268	-2.63	+0.52		0.32	-0.62	0.11					-0.98					(60)
HD 108317	-2.37	+0.48	1.44	0.48	-0.42	0.45	-0.86	-0.05	-0.04	-0.84	-0.67	-1.34	-1.15	0.06		(55, 63–65)
HD 160617	-1.77	+0.44	1.71	1.22	0.25	1.05		0.61	0.59		-0.07	-0.89	-0.03	0.38	0.67	(66, 67)
BD +08°2548	-2.11	+0.43		0.78	-0.19	0.65					-0.36	-0.80				(61, 62)
HD 84937	-2.25	+0.38	1.23	1.04	-0.03	0.65		0.39	0.47			-0.17	-0.44	0.33		(55, 67)
HD 19445	-2.15	+0.37	1.56	1.00	0.13	0.82		0.48	0.44			-0.37	0.11	0.65		(55, 67)
HD 108577	-2.36	+0.36		0.35	-0.56	0.23					-0.69	-1.31				(61, 62)
HD 107752	-2.85	+0.31		-0.25	-0.86	-0.20		-0.90	-0.96		-1.35					(68)
CS 22966-057	-2.62	+0.29		0.20	-0.64	-0.06		-0.15	0.05		-0.70					(49, 50)
HD 126238	-1.93	+0.22		0.98	-0.10	0.74	-0.36	0.18	0.09	-0.65	-0.32	-1.01				(64)
HD 2796	-2.47	+0.17		0.50	-0.48	0.21		-0.40	-0.35		-0.70					(49, 50)
CS 29518-051	-2.69	+0.15		0.41	-0.51	0.16		-0.50	-0.70		-0.90					(49, 50)
HD 85773	-2.62	+0.02		0.01	-0.86	-0.22		-0.97	-1.01		-1.30					(68)
HD 128279	-2.46	-0.02	0.57	-0.18	-0.89	0.00	-0.88	-0.56	-0.63		-1.24		-1.32	-0.12		(55, 63–65)
HD 110184	-2.52	-0.03		0.46	-0.76	-0.03		-0.70	-0.85		-1.22					(68)
CS 22873-055	-2.99	-0.19		-0.42	-1.28	-0.56		-1.00			-1.30					(49, 50)
BD -18°5550	-3.06	-0.22		-1.15	-1.91	-1.17		-1.50			-1.70					(49, 50)
HD 140283	-2.57	-0.22	0.65	0.07	-0.77	-0.05		-0.44				-1.34	-1.06			(67)
CS 22873-166	-2.97	-0.32		-0.05	-0.86	-0.17		-0.70	-1.20		-1.30					(49, 50)
HD 88609	-3.07	-0.33		-0.19	-1.03	-0.20	-1.70	-1.00	-0.91		-1.35	-2.06				(69)
HD 122563	-2.77	-0.52		-0.10	-0.97	-0.25	-1.46	-0.87	-0.86		-1.36	-1.91				(64, 69, 70)

Table S1: Stellar abundances for lighter r-process elements, sorted by decreasing [Eu/Fe]

416

ratios.

418

Star	[Fe/H] error	[Eu/Fe] error	$\log \epsilon$ (Se) error	$\log \epsilon$ (Sr) error	$\log \epsilon$ (Y) error	$\log \epsilon$ (Zr) error	$\log \epsilon$ (Nb) error	$\log \epsilon$ (Mo) error	$\log \epsilon$ (Ru) error	$\log \epsilon$ (Rh) error	$\log \epsilon$ (Pd) error	$\log \epsilon$ (Ag) error	$\log \epsilon$ (Cd) error	$\log \epsilon$ (Sn) error	$\log \epsilon$ (Te) error
CS 31082-001	0.10	0.05		0.10	0.07	0.08	0.12	0.13	0.12	0.12	0.07	0.21			
CS 29497-004	0.12	0.06			0.11	0.08	0.07	0.07	0.14	0.15	0.16	0.18	0.17		
CS 22892-052	0.13	0.05			0.13	0.10	0.12	0.09	0.20	0.10	0.15	0.10	0.11		
J1432-4125	0.18	0.14			0.34	0.10	0.11		0.20	0.09		0.20			
HE 1219-0312	0.16	0.10			0.18	0.15	0.14					0.23			
HD 222925	0.10	0.08	0.22		0.13	0.07	0.08	0.11	0.07	0.11	0.12	0.08	0.13	0.17	0.20
J1538-1804	0.10	0.11			0.12	0.14	0.14			0.28	0.28	0.36	0.30		
CS 31078-018	0.10	0.17			0.09	0.15	0.14			0.12		0.12			
CS 22953-003	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20			
HE 2327-5642	0.10	0.07			0.12	0.09	0.09		0.32			0.32			
CS 22896-154	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20	0.30		
BD +17°3248	0.08	0.10			0.10	0.05	0.14	0.20	0.10	0.10	0.10	0.09	0.15	0.16	0.30
HD 221170	0.12	0.07			0.08	0.07	0.09	0.30	0.10	0.05	0.23	0.05	0.10		
CD -45°3283	0.10	0.17			0.30	0.15	0.17					0.19	0.25		
HD 20	0.04	0.10			0.26	0.09	0.09		0.11	0.11	0.40	0.21	0.21		
HD 120559	0.10	0.17			0.12	0.15	0.17					0.21	0.25		
HD 3567	0.10	0.18			0.20	0.19	0.17					0.20	0.25		
BS 17569-049	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20	0.30		
J1830-4555	0.10	0.20			0.20	0.20	0.25			0.25		0.25			
HD 115444	0.15	0.10			0.11	0.08	0.17					0.11			
HD 186478	0.16	0.08			0.30	0.12	0.06					0.25	0.25		
HD 6268	0.11	0.10			0.33	0.06	0.20					0.17			
HD 108317	0.22	0.18	0.42	0.19	0.20	0.20	0.19	0.28	0.16	0.27	0.18	0.21	0.46		0.30
HD 160617	0.29	0.29	0.21	0.30	0.30	0.30		0.14	0.10		0.22	0.28	0.23	0.10	0.32
BD +08°2548	0.19	0.06			0.30	0.15	0.09				0.25	0.25			
HD 84937	0.10	0.10	0.10	0.10	0.10	0.10		0.10	0.10				0.10	0.18	0.11
HD 19445	0.10	0.10	0.23	0.10	0.10	0.10		0.10	0.16			0.15	0.10	0.15	
HD 108577	0.13	0.12			0.30	0.13	0.10				0.25	0.25			
HD 107752	0.10	0.16			0.28	0.16	0.16		0.17	0.21		0.18			
CS 22966-057	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20			
HD 126238	0.19	0.20			0.24	0.19	0.24	0.20	0.23	0.17	0.38	0.17	0.38		
HD 2796	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20			
CS 29518-051	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20			
HD 85773	0.10	0.20			0.20	0.28	0.19		0.26	0.28		0.22			
HD 128279	0.22	0.18	0.36	0.19	0.21	0.20	0.27	0.28	0.19		0.27		0.40		0.30
HD 110184	0.10	0.16			0.28	0.18	0.21		0.17	0.21		0.18			
CS 22873-055	0.10	0.10			0.20	0.20	0.20		0.20			0.20			
BD -18°5550	0.10	0.10			0.20	0.20	0.20		0.20			0.20			
HD 140283	0.10	0.10	0.10	0.10	0.10	0.10		0.10					0.10	0.10	
CS 22873-166	0.10	0.10			0.20	0.20	0.20		0.20	0.20		0.20			
HD 88609	0.20	0.12			0.12	0.10	0.16	0.12	0.12	0.12		0.12	0.12		
HD 122563	0.19	0.14			0.14	0.09	0.16	0.14	0.14	0.14		0.14	0.14		

419 **Table S2: Uncertainties in the abundances and abundance ratios presented in Table S1.** All
420 values are expressed in dex. In cases where the original study did not present an uncertainty, or
421 where only the statistical error was presented, we adopt a minimum total (statistical plus
422 systematic) uncertainty of 0.10 dex. References are given in Table S1.
423

Star	[Fe/H]	[Eu/Fe]	log ε (Ba)	log ε (La)	log ε (Ce)	log ε (Pr)	log ε (Nd)	log ε (Sm)	log ε (Gd)	log ε (Tb)	log ε (Dy)	log ε (Ho)	log ε (Er)	log ε (Tm)	log ε (Yb)	log ε (Hf)	log ε (Os)	log ε (Pt)	log ε (Th)
CS 31082-001	-2.90	+1.69	0.40	-0.62	-0.29	-0.79	-0.15	-0.42	-0.21	-1.01	-0.07	-0.80	-0.30	-1.15	-0.41	-0.72	0.18	0.30	-0.98
CS 29497-004	-2.84	+1.67	0.35	-0.38	-0.19	-0.67	-0.02	-0.30	-0.31	-1.08	-0.08	-0.76	-0.29	-1.18	-0.22	-0.64	0.41	0.50	-1.16
CS 22892-052	-3.09	+1.64	-0.01	-0.84	-0.46	-0.96	-0.37	-0.61	-0.42	-1.13	-0.26	-0.92	-0.48	-1.39	-0.55	-0.88	0.04	0.20	-1.57
J1432-4125	-2.97	+1.44	0.03	-0.86	-0.53	-0.98	-0.34	-0.67	-0.51	-1.22	-0.39	-1.18	-0.66	-1.45	-0.79	-0.84	0.00		-1.47
HE 1219-0312	-2.97	+1.38	-0.14	-0.86	-0.52	-1.17	-0.41	-0.59	-0.41		-0.34	-1.13	-0.56	-1.51		-0.97			-1.29
HD 222925	-1.46	+1.32	1.26	0.51	0.85	0.22	0.88	0.62	0.82	0.18	1.01	0.12	0.73	-0.09	0.55	0.32	1.17	1.45	-0.06
J1538-1804	-2.09	+1.27	0.69	-0.24	0.08	-0.47	0.17	-0.12	0.06	-0.71	0.31	-0.54	-0.13	-0.83	0.25	-0.36	0.42		-0.97
CS 31078-018	-2.84	+1.23	-0.06	-1.00	-0.67		-0.62	-0.77	-0.76		-0.77	-1.11	-0.96		-1.07				-1.35
CS 22953-003	-2.84	+1.03	-0.22	-1.01	-0.60	-1.45	-0.62	-1.49	-0.71		-0.66	-1.40	-0.85		-0.74				
HE 2327-5642	-2.79	+0.98	-0.30	-1.10	-0.63	-1.15	-0.61	-0.92	-0.72	-1.51	-0.60	-1.31	-0.80	-1.65	-1.03	-1.32	-0.17		-1.67
CS 22896-154	-2.69	+0.96	-0.05	-1.10	-0.40	-1.25	-0.52	-0.90	-0.70		-0.58	-1.55	-0.75						
BD +17°3248	-2.10	+0.90	0.48	-0.42	-0.11	-0.71	-0.09	-0.34	-0.14	-0.91	-0.04	-0.70	-0.25	-1.12	-0.27	-0.57	-0.11	0.50	-1.18
HD 221170	-2.20	+0.80	0.18	-0.73	-0.42	-1.00	-0.35	-0.66	-0.46	-1.21	-0.29	-0.97	-0.47	-1.39	-0.48	-0.84	0.16	0.48	-1.46
CD -45°3283	-0.99	+0.78	1.51																
HD 20	-1.60	+0.73	0.93	-0.09	0.19	-0.35	0.21	-0.06	0.04	-0.74	0.21	-0.49	-0.04	-0.87	-0.06	-0.23	0.40		-0.85
HD 120559	-1.31	+0.71	1.09																
HD 3567	-1.33	+0.70	1.11																
BS 17569-049	-2.88	+0.70	-0.55	-1.33	-1.07	-1.42	-0.95	-0.98	-1.14		-1.15	-1.90	-1.40		-1.20				
J1830-4555	-3.57	+0.69	-1.04	-2.13	-1.57	-2.36	-1.91	-2.16	-2.00		-1.95	-2.60	-2.15	-2.85	-2.46				
HD 115444	-2.85	+0.66	-0.49	-1.53	-1.10	-1.49	-1.06	-1.24	-1.10	-2.29	-1.07		-1.26	-2.18	-1.46				-1.97
HD 186478	-2.60	+0.52	-0.55	-1.38	-1.15		-0.96	-1.30	-1.06		-1.12		-1.15		-1.33				
HD 6268	-2.63	+0.52	-0.45	-1.32	-0.88	-1.54	-0.89	-1.27	-1.01	-2.15	-1.00		-1.20	-2.22	-1.33				-1.93
HD 108317	-2.37	+0.48	-0.32	-1.12	-0.68	-0.90	-0.71	-1.05	-0.82	-1.30	-0.75	-1.50	-0.94	-1.80	-1.13		-0.68	-0.46	-1.99
HD 160617	-1.77	+0.44	0.44	-0.36	0.06		-0.03	-0.45	-0.30		-0.20		-0.35		-0.42				
BD +08°2548	-2.11	+0.43	-0.07	-0.96	-0.65	-1.61	-0.45	-0.90	-0.85	-1.84	-0.46		-0.81	-1.71	-0.82				
HD 84937	-2.25	+0.38	-0.22																
HD 19445	-2.15	+0.37	0.00																
HD 108577	-2.36	+0.36	-0.35	-1.24	-1.05		-0.82	-1.15	-1.11		-0.99		-1.09	-1.98	-1.10				
HD 107752	-2.85	+0.31	-1.09	-1.83	-1.26		-1.23												
CS 22966-057	-2.62	+0.29	-0.73	-1.20		-0.65			-1.00		-1.05								
HD 126238	-1.93	+0.22	-0.03	-0.90	-0.49	-1.08	-0.52	-0.93	-0.68	-1.61	-0.60	-1.37	-0.78	-1.60	-0.93	-0.99	-0.38	-0.34	
HD 2796	-2.47	+0.17	-0.48	-1.40	-0.83	-1.50	-1.10	-1.40	-1.40		-1.32	-2.20	-1.43		-1.52				
CS 29518-051	-2.69	+0.15	-1.01	-2.01		-1.56	-1.18												
HD 85773	-2.62	+0.02	-0.77	-1.75	-1.40		-1.27												
HD 128279	-2.46	-0.02	-1.03	-1.64	-1.15		-1.23	-1.50			-1.35	-1.71	-1.54	-2.24	-1.69				
HD 110184	-2.52	-0.03	-1.06	-1.73	-1.26		-1.30												
CS 22873-055	-2.99	-0.19	-1.31	-2.29		-2.55	-1.78		-1.90		-2.09		-2.30						
BD -18°5550	-3.06	-0.22	-1.67			-2.02		-2.20			-3.00	-2.25							
HD 140283	-2.57	-0.22	-1.09																
CS 22873-166	-2.97	-0.32	-1.54	-2.57	-1.73	-2.30	-1.72					-2.40							
HD 88609	-3.07	-0.33	-1.71	-2.75	-2.02	-2.20	-2.11	-2.41	-2.83		-2.85		-2.79		-2.94				
HD 122563	-2.77	-0.52	-1.65	-2.60	-1.91	-2.15	-2.01	-2.16	-2.44		-2.62		-2.66		-2.78				

424 **Table S3: Stellar abundances for the heavier r-process elements, sorted by decreasing**

[Eu/Fe] ratios. References are given in Table S1.

Star	$\log \epsilon$ (Ba) error	$\log \epsilon$ (La) error	$\log \epsilon$ (Ce) error	$\log \epsilon$ (Pr) error	$\log \epsilon$ (Nd) error	$\log \epsilon$ (Sm) error	$\log \epsilon$ (Gd) error	$\log \epsilon$ (Tb) error	$\log \epsilon$ (Dy) error	$\log \epsilon$ (Ho) error	$\log \epsilon$ (Er) error	$\log \epsilon$ (Tm) error	$\log \epsilon$ (Yb) error	$\log \epsilon$ (Hf) error	$\log \epsilon$ (Os) error	$\log \epsilon$ (Pt) error	$\log \epsilon$ (Th) error
CS 31082-001	0.14	0.05	0.05	0.05	0.05	0.05	0.05	0.05	0.05	0.06	0.05	0.05	0.11	0.05	0.07	0.23	0.13
CS 29497-004	0.12	0.08	0.06	0.06	0.07	0.06	0.07	0.06	0.07	0.17	0.11	0.08	0.20	0.07	0.12	0.21	0.08
CS 22892-052	0.12	0.05	0.05	0.07	0.06	0.07	0.07	0.04	0.06	0.02	0.04	0.04	0.10	0.04	0.15	0.15	0.10
J1432-4125	0.13	0.05	0.08	0.06	0.08	0.05	0.09	0.07	0.06	0.07	0.09	0.14	0.10	0.26	0.10		0.10
HE 1219-0312	0.13	0.12	0.21	0.14	0.12	0.15	0.15			0.13	0.17	0.14	0.16				0.14
HD 222925	0.06	0.07	0.07	0.08	0.08	0.08	0.08	0.09	0.08	0.12	0.08	0.08	0.19	0.10	0.09	0.10	0.11
J1538-1804	0.08	0.12	0.13	0.13	0.12	0.12	0.12	0.12	0.07	0.07	0.07	0.08	0.11	0.13	0.28		0.14
CS 31078-018	0.31	0.22	0.21		0.21	0.16	0.17			0.15	0.22	0.20					0.25
CS 22953-003	0.20	0.20	0.20	0.20	0.20	0.20				0.20	0.20	0.20					0.20
HE 2327-5642	0.03	0.02	0.06	0.06	0.08	0.09	0.07	0.03	0.06	0.10	0.09	0.08	0.10	0.10	0.10		0.10
CS 22896-154	0.20	0.20	0.20	0.20	0.20	0.20				0.20	0.20	0.20					
BD +17°3248	0.11	0.05	0.05	0.06	0.06	0.05	0.04	0.05	0.05	0.05	0.04	0.05	0.01	0.08	0.16	0.15	0.10
HD 221170	0.11	0.06	0.04	0.07	0.08	0.07	0.14	0.08	0.06	0.07	0.08	0.06	0.10	0.11	0.10	0.15	0.05
CD -45°3283	0.15																
HD 20	0.10	0.06	0.03	0.10	0.06	0.04	0.15	0.10	0.07	0.10	0.09	0.10	0.20	0.10	0.10		0.10
HD 120559																	
HD 3567																	
BS 17569-049	0.20	0.20	0.20	0.20	0.20	0.20	0.20			0.20	0.20	0.20					0.20
J1830-4555	0.15	0.20	0.20	0.30	0.20	0.30	0.25			0.20	0.20	0.20	0.25	0.25			
HD 115444	0.11	0.05	0.10	0.11	0.06	0.19	0.21	0.11	0.09		0.09	0.11	0.11				0.15
HD 186478	0.22	0.09	0.11		0.11	0.12	0.15			0.10		0.06		0.27			
HD 6268	0.17	0.05	0.11	0.17	0.08	0.21	0.17	0.17	0.14		0.13	0.07	0.17				0.10
HD 108317	0.19	0.17	0.17	0.23	0.18	0.20	0.19	0.29	0.18	0.20	0.18	0.19	0.22		0.36	0.19	0.21
HD 160617	0.30	0.30	0.30		0.30	0.33	0.30			0.29		0.30		0.33			
BD +08°2548	0.24	0.09	0.11	0.13	0.12	0.11	0.22	0.17	0.12		0.09	0.06	0.30				
HD 84937																	
HD 19445																	
HD 108577	0.21	0.13	0.14		0.14	0.14	0.17			0.11		0.11	0.19	0.28			
HD 107752	0.18	0.16	0.16			0.16											
CS 22966-057	0.20	0.20			0.20					0.20		0.20					
HD 126238	0.21	0.18	0.18	0.21	0.19	0.19	0.19	0.26	0.18	0.40	0.26	0.20	0.26	0.28	0.41	0.42	
HD 2796	0.20	0.20	0.20	0.20	0.20	0.20	0.20		0.20	0.20	0.20		0.20				
CS 29518-051	0.20	0.20		0.20	0.20												
HD 85773	0.18	0.18	0.19		0.19												
HD 128279	0.18	0.19	0.18		0.18	0.24			0.20	0.29	0.19	0.21	0.21				
HD 110184	0.18	0.16	0.16		0.16												
CS 22873-055	0.20	0.20		0.20	0.20		0.20		0.20		0.20		0.20				
BD -18°5550	0.20				0.20		0.20			0.20	0.20						
HD 140283	0.10																
CS 22873-166	0.20	0.20	0.20	0.20	0.20						0.20						
HD 88609	0.10	0.15	0.18	0.12	0.12	0.12	0.12		0.20		0.20		0.12				
HD 122563	0.12	0.16	0.17	0.14	0.16	0.14	0.14		0.14		0.14		0.14				

426 **Table S4: Uncertainties in the abundances and abundance ratios presented in Table S3.** All
427 values are expressed in dex. In cases where the original study did not present an uncertainty, or
428 where only the statistical error was presented, we adopt a minimum total (statistical plus
429 systematic) uncertainty of 0.10 dex. References are given in Table S1.

Element X	$\log \varepsilon (X/Zr)_{\text{base}}$	standard error	n
Se	0.64	0.10	2
Sr	0.16	0.04	13
Y	-0.73	0.02	13
Zr	0.00	0.05	13
Nb	-1.17	0.11	4
Mo	-0.54	0.05	13
Ru	-0.66	0.09	10
Rh	-1.39	0.10	1
Pd	-0.99	0.06	12
Ag	-1.76	0.05	3
Cd	-1.31	0.10	2
Sn	-1.01	0.10	1
Te	-0.12	0.10	1

430 **Table S5: The empirical baseline abundance pattern for lighter elements calculated from**
 431 **the abundances of the 13 stars with $[\text{Eu}/\text{Fe}] \leq +0.3$.** We adopt a minimum uncertainty of 0.10
 432 dex when the number of stars, n , used to compute the mean is ≤ 2 .

Element X	$\log \varepsilon (X/\text{Ba})_{\text{base}}$	standard error	n
Ba	0.00	0.05	13
La	-0.87	0.06	11
Ce	-0.32	0.05	8
Pr	-0.80	0.11	7
Nd	-0.33	0.05	12
Sm	-0.70	0.08	5
Eu	-1.15	0.03	13
Gd	-0.77	0.08	6
Tb	-1.58	0.10	1
Dy	-0.70	0.11	7
Ho	-1.27	0.19	4
Er	-0.78	0.08	9
Tm	-1.39	0.13	2
Yb	-0.99	0.09	5
Hf	-0.96	0.10	1
Os	-0.35	0.10	1
Pt	-0.31	0.10	1

433 **Table S6: The empirical baseline abundance pattern for heavier elements calculated from**
 434 **the abundances of the 13 stars with $[\text{Eu}/\text{Fe}] \leq +0.3$.** We adopt a minimum uncertainty of 0.10
 435 dex when the number of stars, n , used to compute the mean is ≤ 2 .